

## LECTURE # 3

•  $\chi f \leftrightarrow \chi f$  IN EARLY UNIVERSE

• INTRO TO IND. DET.

[IF TIME PERMITS...]

THE DM THERMAL HISTORY IS POTENTIALLY IMPORTANT FOR DM STRUCTURE FORMATION, ESP FOR (COLD) "WIMPS"

IN EARLY UNIVERSE,  $\chi A \leftrightarrow \chi f$  ( $f \in$  THERMAL BATH)

KEEPS DM IN KINETIC EQ EVEN AFTER CHEMICAL FREEZE-OUT (i.e. WHEN  $\Gamma_{\chi\chi\leftrightarrow\chi} \ll H$ )

REASON IS THAT TARGET DENSITIES ARE VERY DIFF:

$$\chi A \leftrightarrow f f \rightarrow \Gamma = M_{NR} \sigma \quad \text{--- EXP. SUPPRESSED}$$

$$\chi f \leftrightarrow \chi f \rightarrow \Gamma = M_{REL} \sigma$$

Typically, FOR WIMP DM  $\sigma_{\chi f \rightarrow \chi f} \sim G_F^2 T^2$

HERE, TO ESTIMATE KINETIC DECOUPLING, WE NEED TO ACCOUNT FOR "INEFFICIENT" MOMENTUM TRANSFER:

$\delta p \sim T$   
TYPICAL MOMENTUM TRANSFER

BUT FOR COLD RELICS  $\frac{p^2}{2M} \sim T$   
IN KINETIC EQUILIBRIUM

SO  $p \sim \sqrt{mT} \gg \delta p$

MOM. TRANSFER IS STOCHASTIC, SO IT TAKES  $\frac{1}{N} \left(\frac{\delta p}{p}\right)^2 \sim \frac{T^2}{mT}$   
SCATTERING TO ESTABLISH KINETIC EQUILIBRIUM  $\frac{T}{m}$

So, here, we want to compare

$$T^3 \cdot G_P^2 T^2 \cdot \frac{T}{M_N} \sim \frac{T^2}{M_P}$$

$\underbrace{\hspace{1.5cm}}_{M_{REL}} \quad \underbrace{\hspace{1.5cm}}_{D_{KINETIC}} \quad \underbrace{\hspace{1.5cm}}_{\left(\frac{S_P}{P}\right)^2} \quad \sim \quad \frac{T^2}{M_P}$

$$So \quad T_{kd} \sim \left( \frac{M_N}{M_P G_P^2} \right)^{1/4} \sim \left( \frac{100 \text{ GeV}}{10^{18} \text{ GeV} \cdot 10^{-10} \text{ GeV}^{-4}} \right)^{1/4} \sim 10^{-6/4} \cdot \left( \frac{M_N}{100 \text{ GeV}} \right)^{1/4} \text{ GeV}$$

$$OR \quad T_{kd} \sim 30 \text{ MeV} \left( \frac{M_N}{100 \text{ GeV}} \right)^{1/4}$$

WHAT DOES THIS IMPLY FOR STRUCTURE FORMATION ?

ROUGHLY, CUTOFF WILL BE SIZE OF HORIZON AT KINETIC DECOUPLING, SO

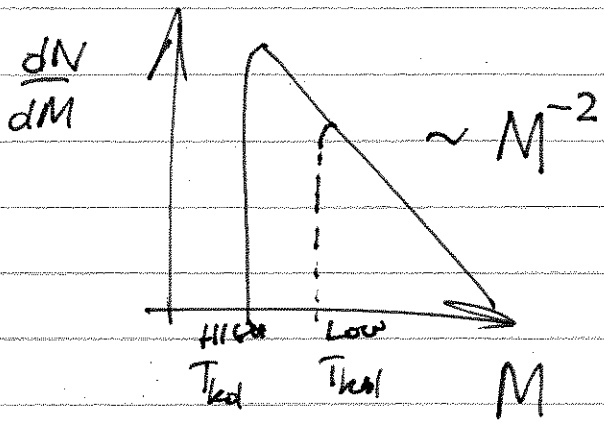
$$M_{hol} \sim \frac{4\pi}{3} \left( \frac{1}{H(T_{kd})} \right)^3 \rho_{DM}(T_{kd})$$

$$\sim 30 M_{\oplus} \left( \frac{10 \text{ MeV}}{T_{kd}} \right)^3$$

(MORE PRECISELY: CUTOFF SCALE SET BY LARGEST OF FREE STREAMING ( $T_{kd} \gtrsim 0.1 \text{ GeV}$ ) VS ACOUSTIC DAMPING ( $T_{kd} \lesssim 0.1 \text{ GeV}$ ))

SO, FOR TYPICAL WIMPS, PROTOHALOS HAVE MASS  $\sim M_\odot$   
( $\rho \sim 10^{-6} M_\odot$ )

IMPORTANT FOR INDIRECT DETECTION (BOOST FACTOR)  
AND FOR DM "SMALL SCALE PROBLEMS"

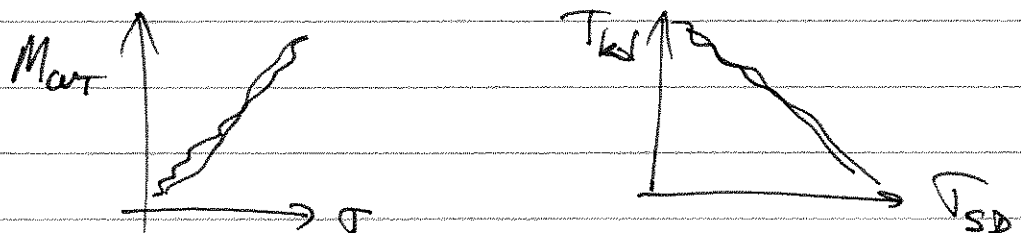


WAYS TO PROBE?  $\chi q \leftrightarrow \chi q$  SAME PROCESS AS DIRECT DETECTION!

$M_{cut}$  CAN CORRELATE WITH  $\sigma_{SI} \rightarrow \sigma_{SI}$

CAVEAT: USUALLY  $T_{kid} < T_{QCD}$ , SO QUARK SCATTERING SECONDARY... ( $\sim 200 \text{ MeV}$ )

BUT IF (MODEL-DEP.) QUARK-LEPTON UNIVERSALITY HOLDS, THEN CORRELATION IS EXPECTED. AND FOUND (CORNELL + PROFUMO, 2012) e.g.: SUSY, UED



# INDIRECT DETECTION

KEY PROCESS:  $\chi\chi \rightarrow SM$  IN LATE UNIVERSE  
(ALSO  $\chi \rightarrow SM$  DECAY)

- KEY INGREDIENTS:
- (i) PRODUCTION RATES
  - (ii) ENERGY SOURCE
  - (iii) ANNIHILATION PRODUCTS

RATE FOR GIVEN PARTICLE, IN UNIT VOLUME

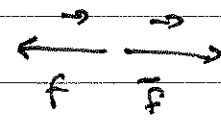
e.g.:  $e^\pm, \bar{p}, \bar{D}, j's, \nu's$

$$= \underbrace{\left( \frac{\# DM PAIRS}{IN V} \right)}_{\int \frac{\rho_{DM}}{m_\chi^2} dV} \times \underbrace{\left( \text{PAIR ANNIHILATION RATE} \right)}_{\langle \sigma v \rangle (T_{=0})} \times \underbrace{\left( \text{\# OF PARTICLES PER ANN. EVENT} \right)}_{m_\chi, \text{ ANN FINAL STATE}}$$


TALKED ABOUT  $\langle \sigma v \rangle, m_\chi$ ; ANN FINAL STATE?

# ANNIHILATION FINAL STATE

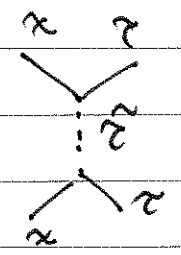
IF DM IS MAJORANA  $\rightarrow$   $\chi\chi \rightarrow f\bar{f}$  REQUIRES  
 HELICITY FLIP  $\rightarrow |M|^2 \propto m_f^2$  (EXACTLY AS FOR  $\pi^0$  DECAY!)

So: NO LIGHT FERMIONS!  
 or 

IF  $m_\chi \lesssim m_{top}$  AND  $\chi$  DOESN'T LIKE  $W^+W^-, \tau\tau, hh$

$\Rightarrow \chi\chi \rightarrow \bar{b}b, \tau^+\tau^-$  DOMINANT  
  
 3X COLOR

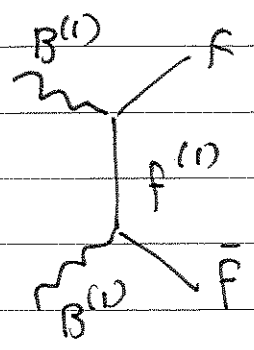
$\tau^+\tau^-$  CAN BE BOOSTED, e.g., WITH LIGHT  $\tilde{\tau}$  IN SUSY.



(TRUE IN MSUGRA/CMSM)

... BUT KILLED BY  $m_h \sim 125$  GeV

TOTALLY OPPOSITE FOR UED,  $B^{(1)} \propto$  HYPERCHARGE



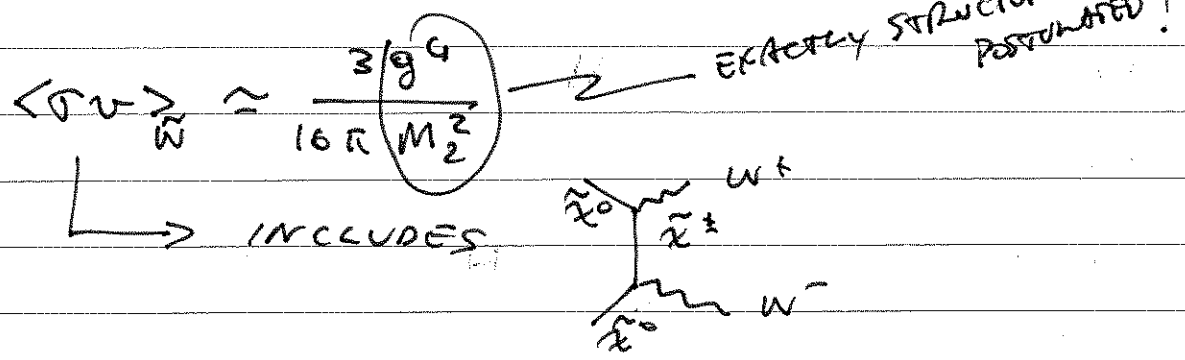
$$|M|^2 \sim |Y_f|^4$$

LIKES  $U_L: |Y_f| = \frac{4}{3}$   
 $e_R: |Y_f| = 2$

IN SU(2) MULTIPLICETS (e.g. HIGGSINOS, WINOS IN SUSY) EVERYTHING IS FIXED BY GAUGE INTERACTIONS

WINOS : FINAL STATES:  $W^+W^-$

(SU(2) TRIPLET)  $\tilde{\chi}^\pm$   
 $\tilde{\chi}^0$   $\approx$   $3 \delta M$   $\sim M_2$   
 $\approx 10^5 \text{ MeV}$



$$\Omega_{\tilde{W}} h^2 \approx 0.1 \left( \frac{M_2}{2.2 \text{ TeV}} \right)^2$$

HIGGSINOS : FINAL STATES:  $W^+W^-, ZZ$

(SU(2) DOUBLET)  $\tilde{\chi}_2^\pm$   
 $\tilde{\chi}_1^0$   $\approx$   $\mu$

$$\langle \sigma v \rangle_H \approx \frac{g^4}{512 \pi \mu^2} (21 + 3 \tan^2 \beta_w + 11 \tan^4 \beta_w)$$

$$\Omega_H h^2 \approx 0.1 \left( \frac{\mu}{1 \text{ TeV}} \right)^2$$

